The Definition of Holography

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Informal Talk

Based on arXiv:2404.00070 [hep-th] and upcoming work with Aayush Verma

June, 2024

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See also more informal notes on AdS/CFT at https://vkalvakotamath.github.io/teaching/2023-ads-cft-1.

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One of the most important revolutions in theoretical high energy physics is that of **holography**, or the duality between a bulk gravitating region and a non-gravitating QFT, usually satisfying conformal invariance.

String theory gives us a way to establish such a duality in the case of anti-de Sitter space, where the bulk is the gravitating AdS and the dual field theory lives on the boundary of AdS. The original description was the duality between type IIB superstring theory and an $\mathcal{N} = 4$ SYM theory, but after that a lot of generality has been found.

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We will try to point out the most general definition of holography and some more technical details in this talk. In the first half, we will discuss what holography generally means and how to approach this phenomenon, and in the second half, we will discuss some more details of a mathematical background to holography.

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We will generally either stick between anti-de Sitter space or de Sitter space holographies.

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Strongly, from bulk reconstruction of bulk fields in the asymptotic boundary limit $r \to \infty$ [from the **extrapolate dictionary** and **HKLL**], we can see that there are only certain bulk fields that can be fully reconstructed given access to a particular boundary subregion. This bulk subregion is called the *entanglement wedge* $\mathcal{E}_W(R)$ for a boundary subregion R.

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Holographic entanglement entropy first indicated this, with Ryu and Takayanagi's famous derivation of the RT formula:

$$S_R = {{\rm Area \ of \ } \gamma_R \over 4G_N \hbar} \ .$$
 (1)

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Note that usually the bulk is the side we are more ignorant about. But here, we are trying to understand how these subregions imply deeper mathematics.

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In general, these are referred to as **functorial QFTs**. Locally covariant QFTs are functorial QFTs that are functors from the category of globally hyperbolic manifolds to the category of C^* -algebras, satisfying some conditions like *causality*, *isotony*, *covariance*, etc.

Generally, when working with LCQFTs, we are concerned with local nets of observables $\{\mathcal{O}_n\}$ associated to a bounded subregion \mathcal{U} . Usual QFTs have a peculiar property, that for such \mathcal{U} on a Cauchy slice Σ , the entanglement entropy S, computed by the area of $\partial \mathcal{U}$, is divergent.

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Without going into the details, the idea is basically that there are three types of von Neumann algebras (cf. Gell Mann Murray and von Neumann), loosely based on the tracial properties that are of concern to us.

To be a little more precise, we consider AQFTs as a functor from the category of $D \ge 2$ globally hyperbolic manifolds Man_D to the category of unital C^* algebras Alg. The objects typically are composed of smooth D-manifolds, with orientation of nowhere vanishing *n*-forms and time-orientation. The morphisms are smooth isometric embeddings that preserve these orientations.

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For general subregions with no asymptotic boundary structures, this von Neumann algebra is typically type III₁ with reduction to type II₁ Jensen, Speranza and Sorce.

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Naively there are three types of von Neumann algebras, but there are sub-classifications based on some intrinsic properties of their algebras.

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Bulk QFTs are no exception, and in the Ryu-Takayanagi formula this is rather explicitly clear.

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This is due to a theorem in mathematics by Takesaki:

$$III_1 \rtimes \mathcal{R}_{aut} = II_{\infty} .$$
 (2)

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III₁ is the usual large N Liu-Leutheusser algebra. \mathcal{R}_{aut} is the modular group of automorphisms of this algebra, and is here the group of time translations generated by the modular operator Δ_{Ψ} . II_{∞} is the resulting algebra, which is an extended large N algebra *including* the Hamiltonian.

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What happens now?

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Important thing: we don't have to rely on Euclidean path integrals! (Path integrals are complicated...)

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By the cosmic No-Hair theorem, almost every class of cosmological models become asymptotically de Sitter. In the context of quantum gravity, this has many interesting consequences, since we can have a general future description of information (see Suvrat et al's Hilbert space work!) Kind of problematic in string theory: vacuas, Maldacena-Nunez, etc. But exciting nonetheless!

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Spoiler alert: It isn't fun.

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Overall, de Sitter has many subtleties, due to which we don't really know precisely what "the right" notion of holography is. Which brings us back to the origianl question: what *is* the most general definition of holography?

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One: duality between a gravitating region and a precisely defined dual field theory living on the "boundary". Operator algebras are also dual to each other, and correspond to subregions on both sides. That is,

$$\mathcal{H}_{\mathsf{Fock}} = \mathcal{H}_{\mathsf{QFT}} , \qquad (3)$$

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Two: Similar idea, but holography of *information*, which is a little different. Usually involves violating the split property, but in general lets us have a duality of operators at asymptotics. So far the de Sitter picture is not clear in either case, but there are many things to look out for.

Thank you for your attention!